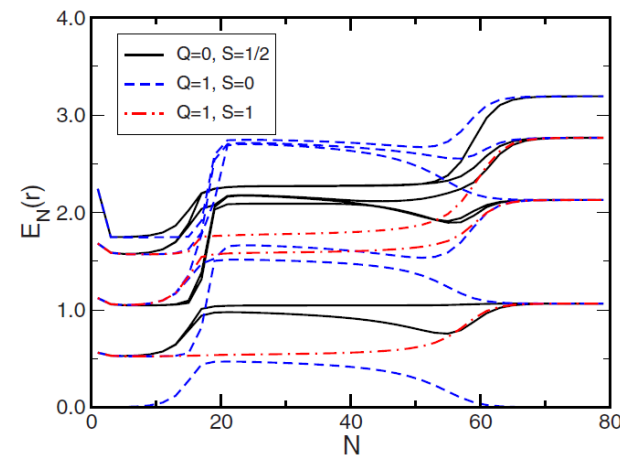
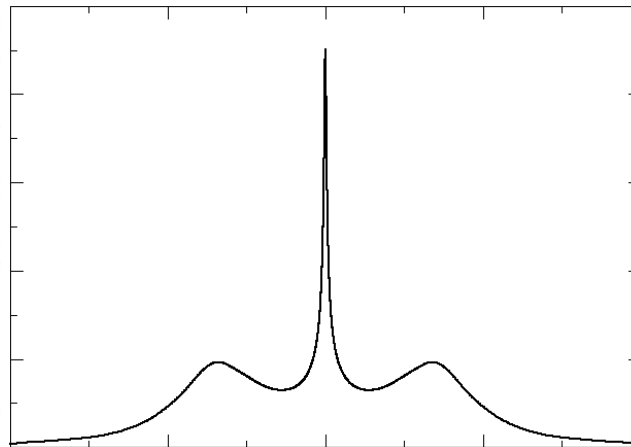
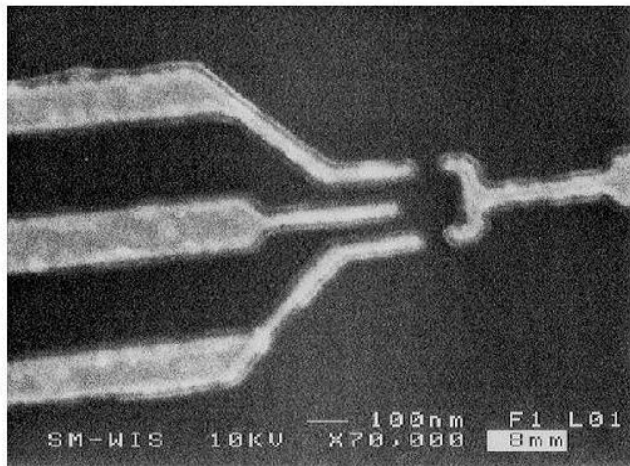


# NUMERICAL METHODS FOR QUANTUM IMPURITY MODELS



<http://www.staff.science.uu.nl/~mitch003/nrg.html>

March 2015

Andrew Mitchell, Utrecht University

# Quantum impurity problems

- **Part 1: Quantum impurity problems and theoretical background**
- **Part 2: Kondo effect and RG. 1d chain formulation and iterative diagonalization**
- **Part 3: Logarithmic discretization and truncation. The RG in NRG**
- **Part 4: Physical quantities. Results and discussion.**

# **NUMERICAL METHODS FOR QUANTUM IMPURITY MODELS**

## **Part 2: Kondo effect and the Renormalization Group**

**March 2015**

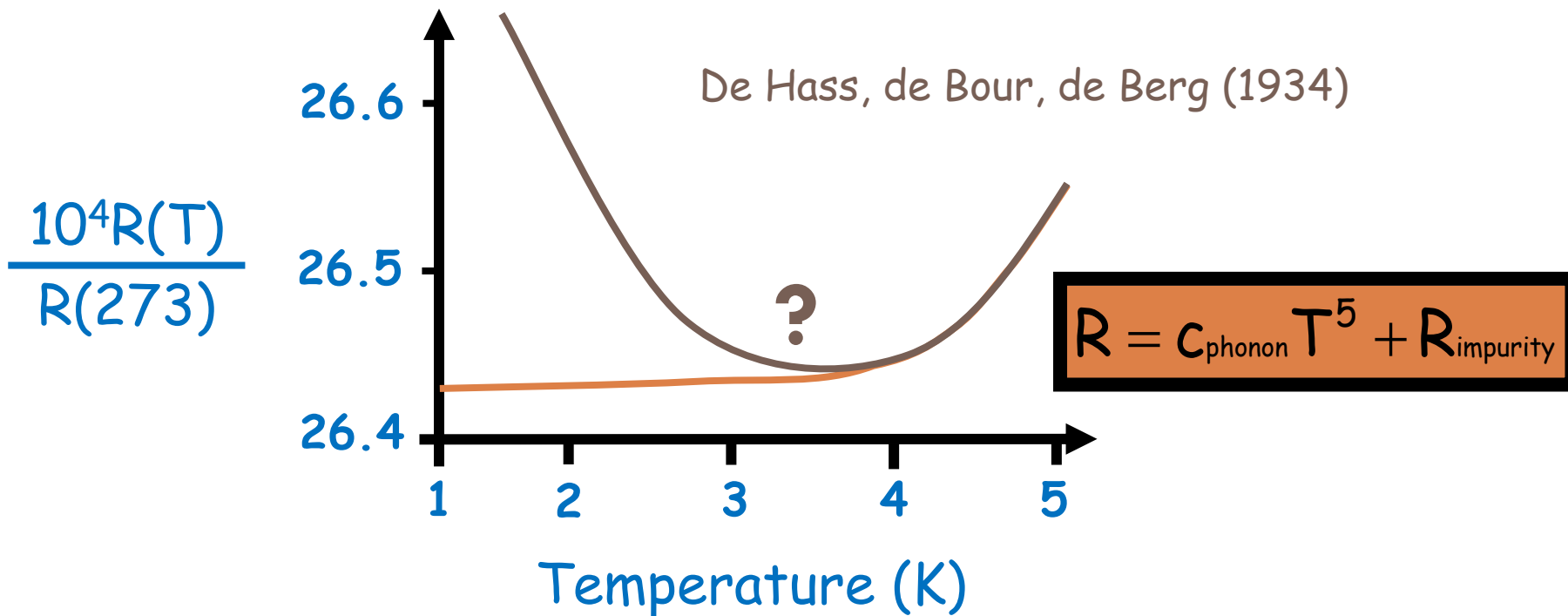
**Andrew Mitchell, Utrecht University**

# Overview: Part 2

- **Kondo effect**
- **Perturbation theory**
- **Perturbative scaling**
  
- **Mapping to 1d chain**
  
- **Iterative diagonalization**

# Magnetic impurities

- Resistance of metals:
- Experiments reveal low-temperature minimum



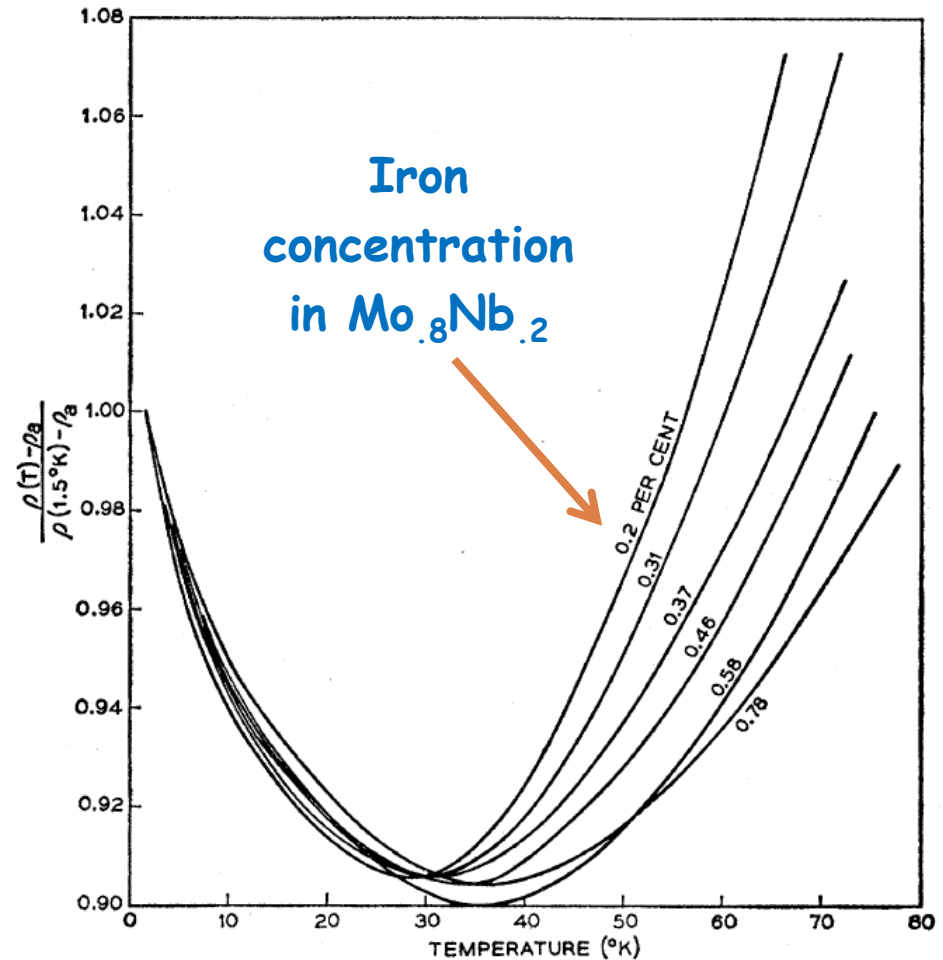
# Magnetic impurities

J. Kondo,  
Prog. Theor. Phys.  
32, 37 (1964)

- Resistance minimum an *impurity* effect

$$T_{\min} \propto (C_{\text{imp}})^{1/5}$$

- Conduction electrons scatter off impurities



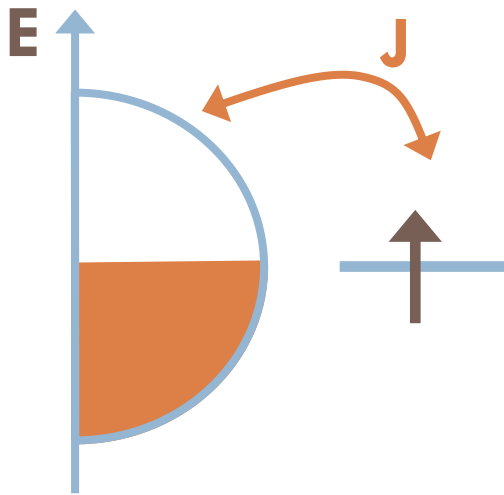
M. Sarachik et al. 1964

# Recap: Kondo model

J. Kondo,  
Prog. Theor. Phys.  
32, 37 (1964)

## □ Scattering from **magnetic** impurities

- Single spin- $1/2$  impurity
- Bath of non-interacting conduction electrons
- AF Exchange coupling



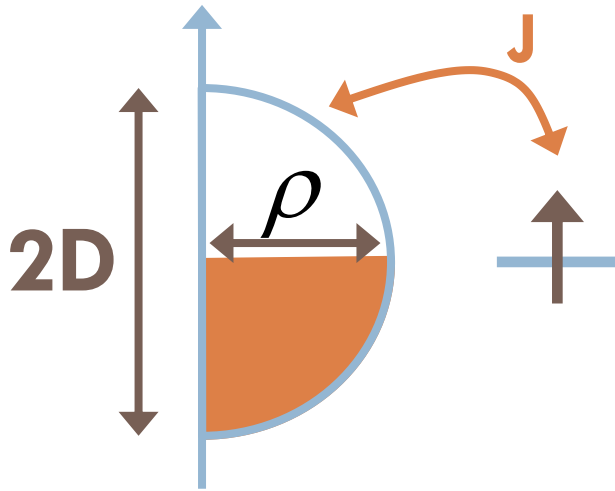
$$H_K = H_{host} + J \vec{S}_{imp} \cdot \vec{S}_0$$
$$\sum_{k\sigma} \varepsilon_k c_{k\sigma}^\dagger c_{k\sigma}$$

impurity  
spin- $1/2$

conduction electron  
spin density

# The Kondo problem

- 3<sup>rd</sup> order perturbation theory in  $J$ :



$$R = C_{\text{phonons}} T^5 + R_{\text{impurity}}$$

$$+ C_{\text{imp}} \left[ J^2 + \rho J^3 \text{Log} \left( \frac{D}{T} \right) \right]$$

Resistance minimum



$$T_{\text{min}} \propto (C_{\text{imp}})^{1/5}$$





# The Kondo problem

J. Kondo,  
Prog. Theor. Phys.  
32, 37 (1964)

- Obviously: perturbation theory **fails** at low T:

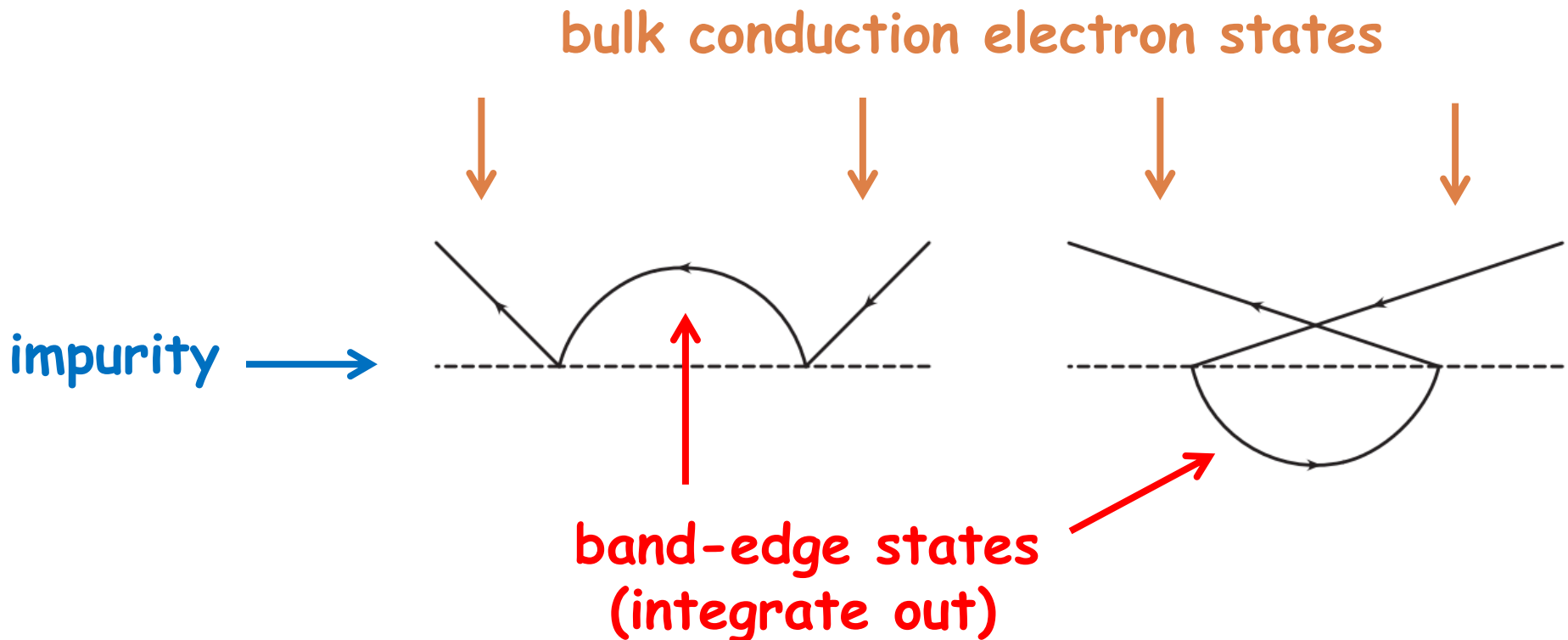
$$\rho J^3 \text{Log}(D/T) > J^2 \quad \rightarrow \quad T_K = D e^{-1/\rho J}$$

- What happens below  $T_K$ ?
- What is the ground state?

# Scaling and Renormalization

- **Look at physics at a lower energy scale**

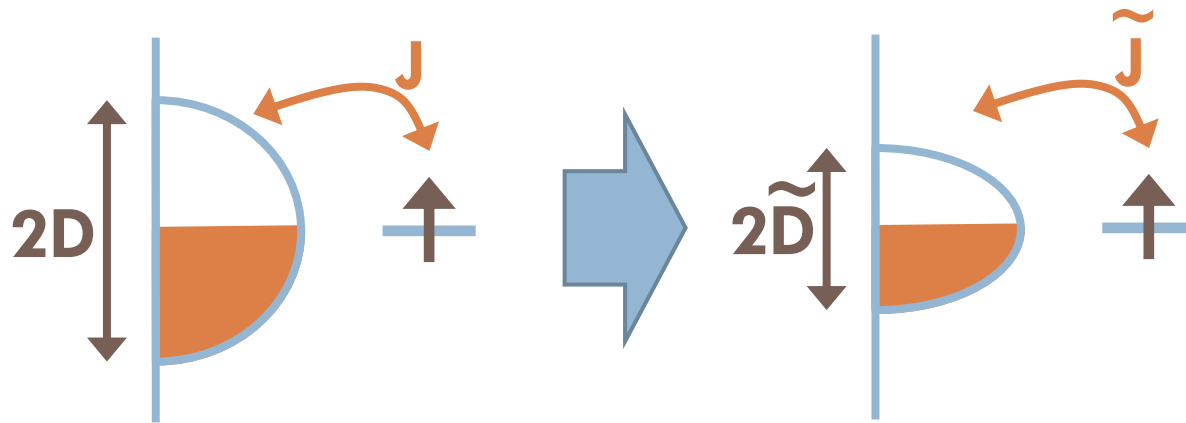
P. W. Anderson, *J. Phys. C* 3, 2436 (1970)



# Scaling and Renormalization

- **Map to a Kondo model of the same form**

P. W. Anderson, *J. Phys. C* 3, 2436 (1970)



Consider a **sequence** of Hamiltonians with different effective (**renormalized**) parameters

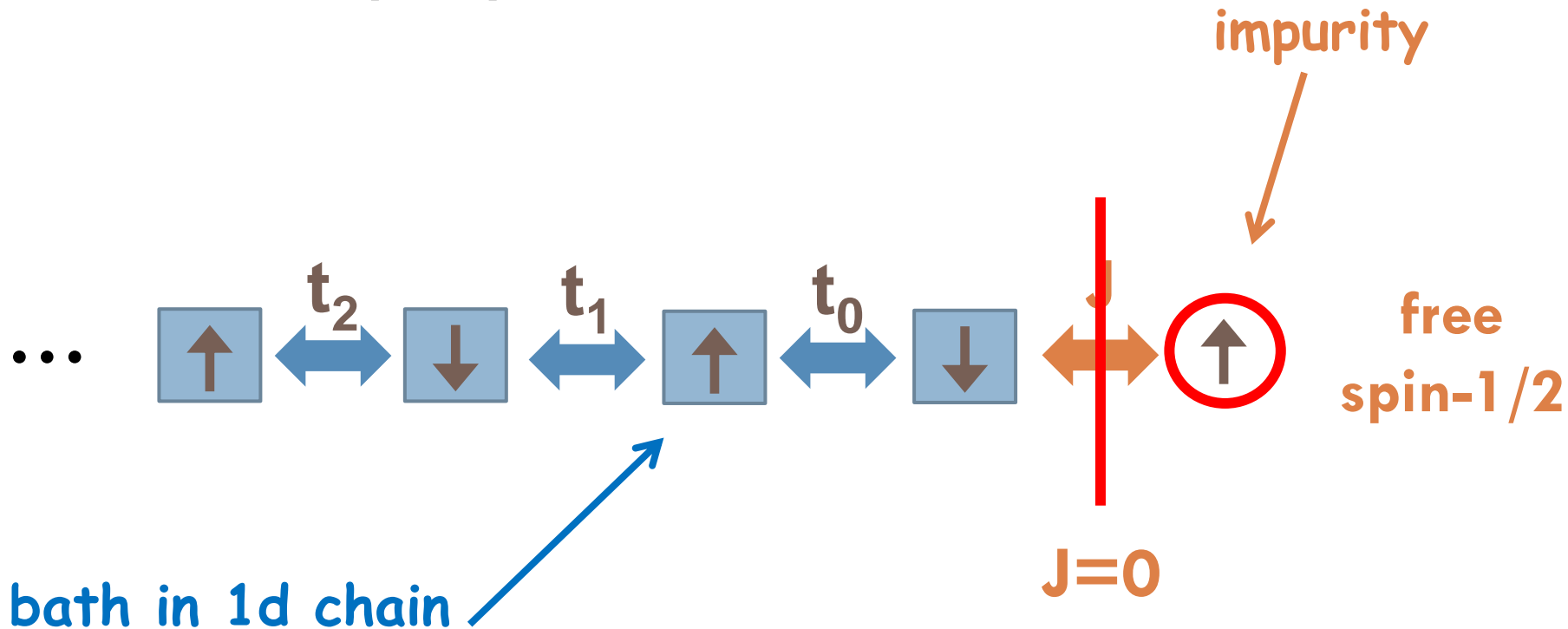
# Fixed points

- **Special values of the parameters produce **NO renormalization** → “fixed points”**
- **Kondo model:**
  - **Local Moment (LM) fixed point:  $J=0$**
  - **Strong coupling (SC) fixed point:  $J=∞$**

# RG fixed points

see Alex Hewson's book  
"The Kondo Problem..."  
CUP (1997)

- **Small J: weak coupling**
  - **free impurity local moment**

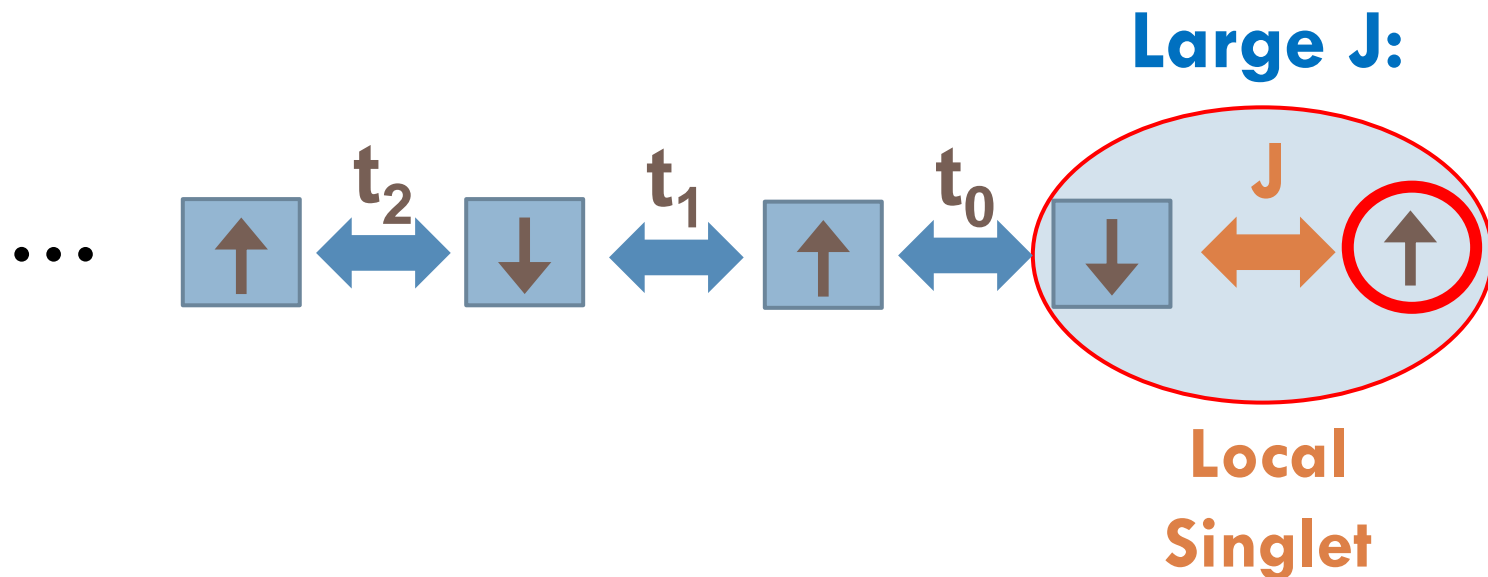


bath in 1d chain  
representation... more on this later!

# RG fixed points

see Alex Hewson's book  
"The Kondo Problem..."  
CUP (1997)

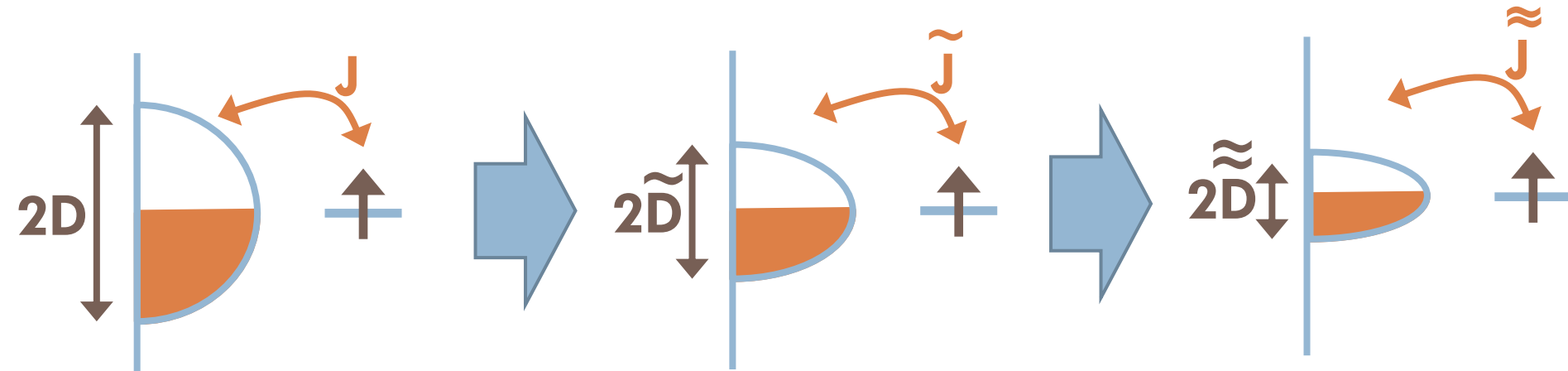
- Large  $J$ : **strong coupling** state
  - Impurity forms **spin-singlet** with conduction electrons



# Scaling and Renormalization

- Look at physics on **successively** lower energy scales

P. W. Anderson, J. Phys. C 3, 2436 (1970)



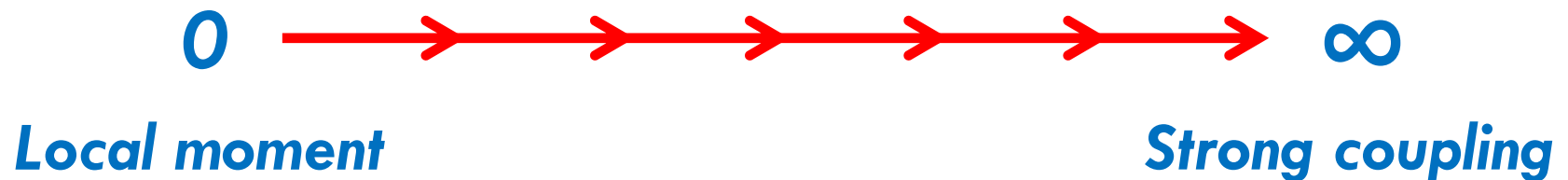
$$\frac{d \rho J}{d \ln D} = -2(\rho J)^2 + \dots$$



$$D e^{-1/\rho J} = \tilde{D} e^{-1/\rho \tilde{J}} \sim k_B T_K$$

# Scaling and Renormalization

- For (initially) weak coupling, perturbative scaling indicates that coupling  $J$  grows under RG

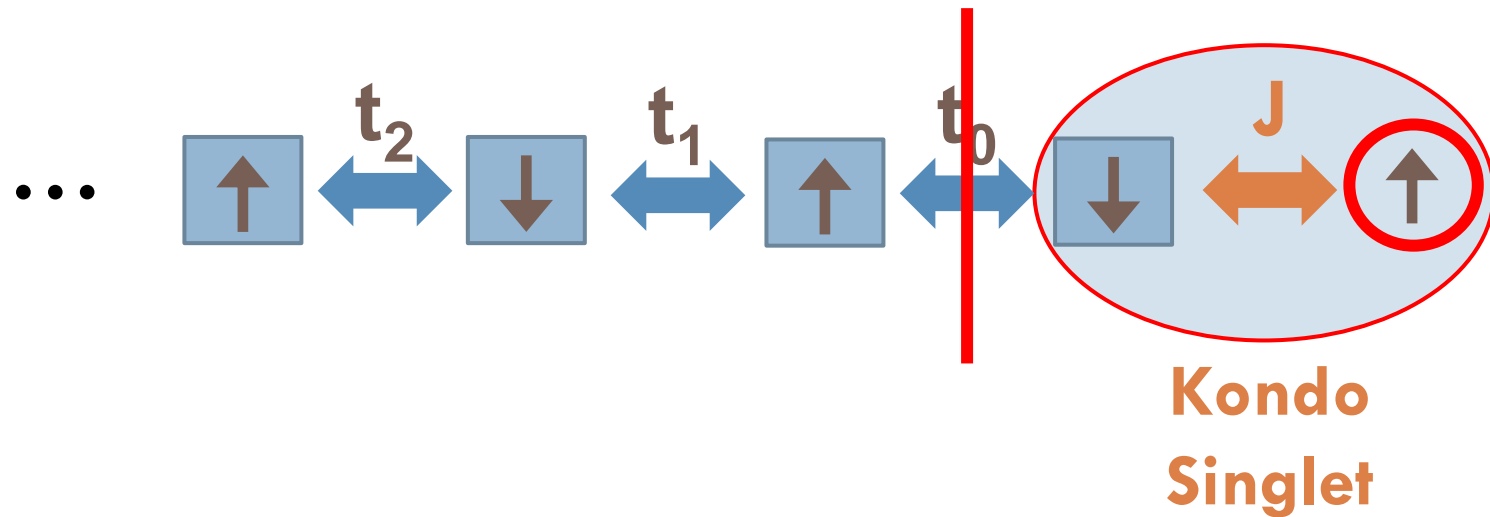




# RG flow

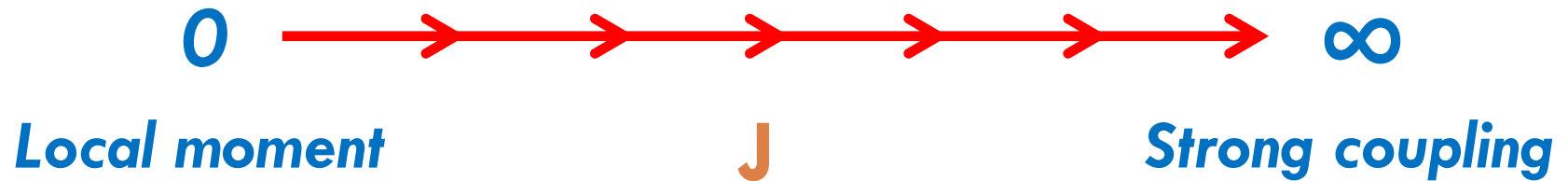
see Alex Hewson's book  
"The Kondo Problem..."  
CUP (1997)

- So: when coupling is initially small...  
... it grows under RG and becomes large



many-body singlet ground state  
(complicated real-space structure)

# RG flow



**BUT: analysis breaks down before  $J \sim O(1)$**

- **Need a **non-perturbative** approach!**
  - ▣ **Must be able to handle *large* energy scales:  $D, J$  etc.**
  - ▣ **... and *exponentially small* scales,  $T_K$**
  - ▣ **Exploit **RG** character of the problem**

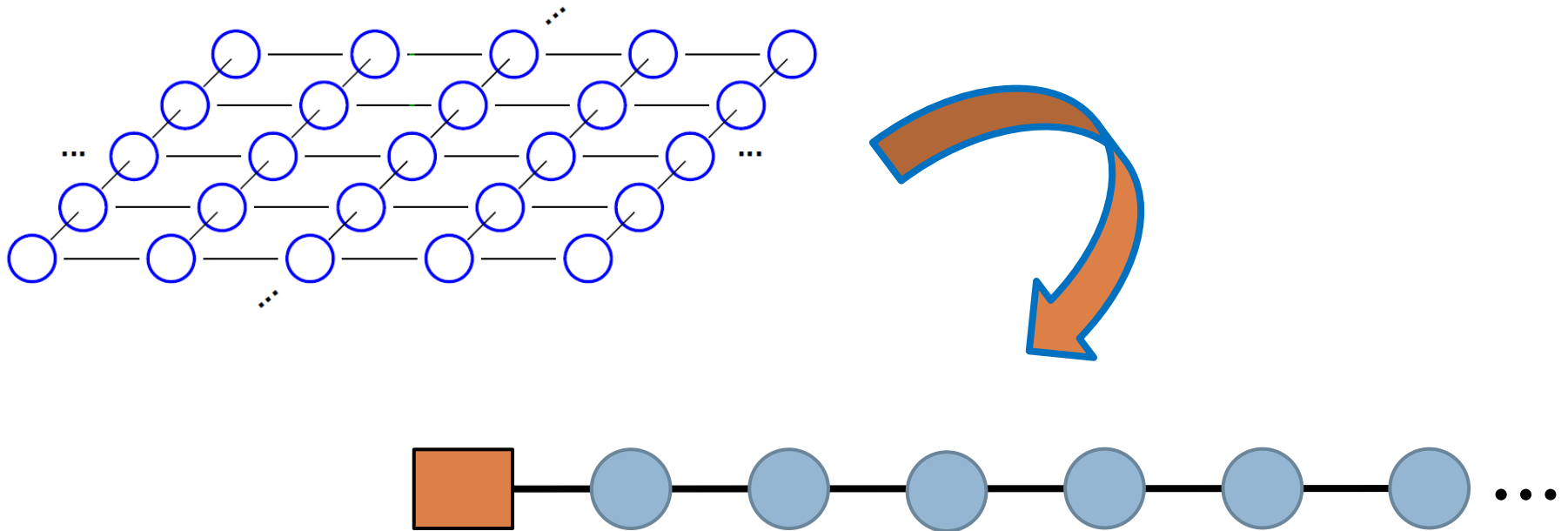
# Numerical Renormalization Group

- **(very) brief description of NRG:**
  - ▣ **Logarithmic discretization of bath**
  - ▣ **Mapping to 1d chain**
  - ▣ **Iterative diagonalization**
  - ▣ **Successive Hilbert-space truncation**

# NRG: preliminaries

## □ 1d chain representation

- Any non-interacting system can be mapped to a **1d tight-binding chain: “Tridiagonalization”**



# Recap: real-space representation

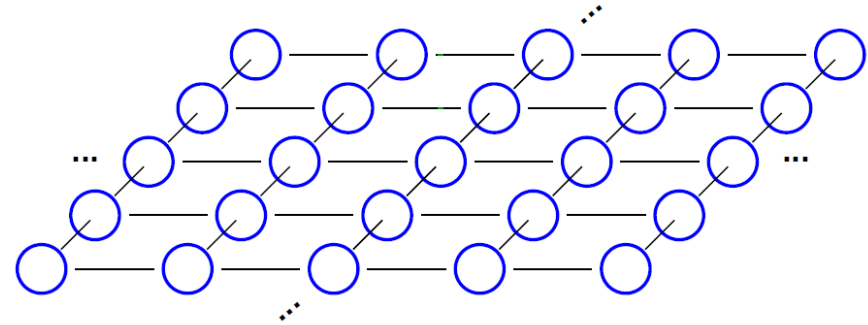
## □ Host metal: non-interacting tight-binding model

$$H_{host} = \sum_{\sigma} \sum_{\langle i,j \rangle} t_{ij} c_{i\sigma}^{\dagger} c_{j\sigma}$$

$$\equiv \sum_{\sigma} \vec{c}_{\sigma}^{\dagger} \underline{T} \vec{c}_{\sigma}$$

$$\vec{c}_{\sigma}^{\dagger} = \left( c_{1\sigma}^{\dagger}, c_{2\sigma}^{\dagger}, c_{3\sigma}^{\dagger}, \dots \right)$$

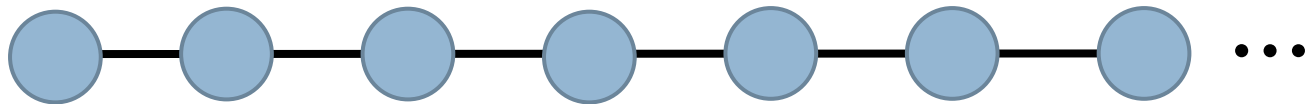
$$; \quad \underline{T} = \begin{bmatrix} t_{11} & t_{12} & \dots \\ t_{12}^* & t_{22} & \\ \vdots & & \ddots \end{bmatrix}$$



# Tridiagonal

□ Host metal: **1d chain** representation

$$\begin{aligned} H_{host} &= \sum_{\sigma} \vec{c}_{\sigma}^{\dagger} \underline{T} \vec{c}_{\sigma} = \sum_{\sigma} \vec{f}_{\sigma}^{\dagger} \underline{W} \vec{f}_{\sigma} \\ &= \sum_{\sigma} \sum_n e_n f_{n\sigma}^{\dagger} f_{n\sigma} + \left( h_n f_{n\sigma}^{\dagger} f_{(n+1)\sigma} + \text{H.c.} \right) \end{aligned}$$



where,  $\vec{f}_{\sigma} = \underline{S} \vec{c}_{\sigma}$  such that,  $\underline{W} = \underline{S}^{\dagger} \underline{T} \underline{S}$

# Tridiagonal

$$\begin{aligned}
 H_{host} &= \sum_{\sigma} \vec{c}_{\sigma}^{\dagger} \underline{\underline{T}} \vec{c}_{\sigma} = \sum_{\sigma} \vec{f}_{\sigma}^{\dagger} \underline{\underline{W}} \vec{f}_{\sigma} \\
 &= \sum_{\sigma} \sum_n e_n f_{n\sigma}^{\dagger} f_{n\sigma} + \left( h_n f_{n\sigma}^{\dagger} f_{(n+1)\sigma} + \text{H.c.} \right)
 \end{aligned}$$

$$\begin{bmatrix}
 t_{11} & t_{12} & t_{13} & t_{14} \\
 t_{21} & t_{22} & t_{23} & t_{24} \\
 t_{31} & t_{32} & t_{33} & t_{34} \\
 t_{41} & t_{42} & t_{43} & t_{44}
 \end{bmatrix}$$

T

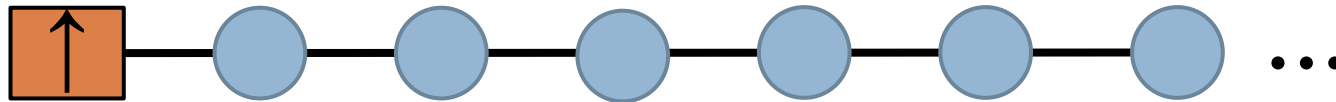
$\Rightarrow$

$$\begin{bmatrix}
 e_1 & h_1 & 0 & 0 \\
 h_1 & e_2 & h_2 & 0 \\
 0 & h_2 & e_3 & h_3 \\
 0 & 0 & h_3 & e_4
 \end{bmatrix}$$

W

# Impurity problem as a 1d chain

## □ 1d chain representation of impurity problem



impurity

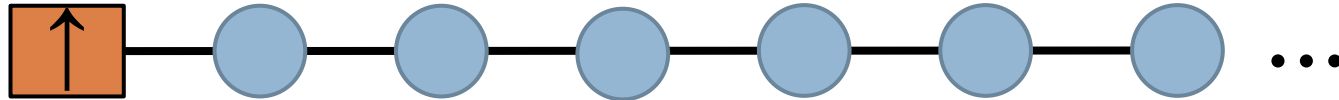
Choose **'zero-orbital'** of the conduction electron chain to be the physical host orbital coupled to the impurity in **real-space**

$$f_{0\sigma} = c_{0\sigma} \quad \Rightarrow \quad \vec{s}_0 = \sum_{\sigma, \sigma'} f_{0\sigma}^\dagger \frac{\vec{\sigma}_{\sigma\sigma'}}{2} f_{0\sigma'}$$



# Impurity problem as a 1d chain

## □ 1d chain representation of impurity problem

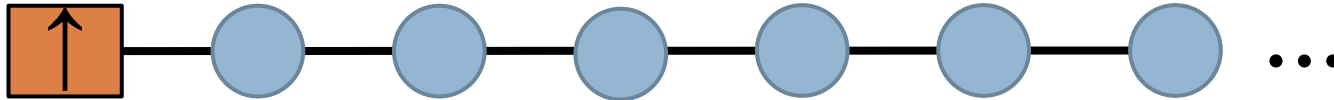


$$H_{Kondo} = H_{host} + H_{imp}$$

$$H_{imp} = J \vec{S}_{imp} \cdot \vec{S}_0$$

$$H_{host} = \sum_{\sigma} \sum_n e_n f_{n\sigma}^{\dagger} f_{n\sigma} + \left( h_n f_{n\sigma}^{\dagger} f_{(n+1)\sigma} + \text{H.c.} \right)$$

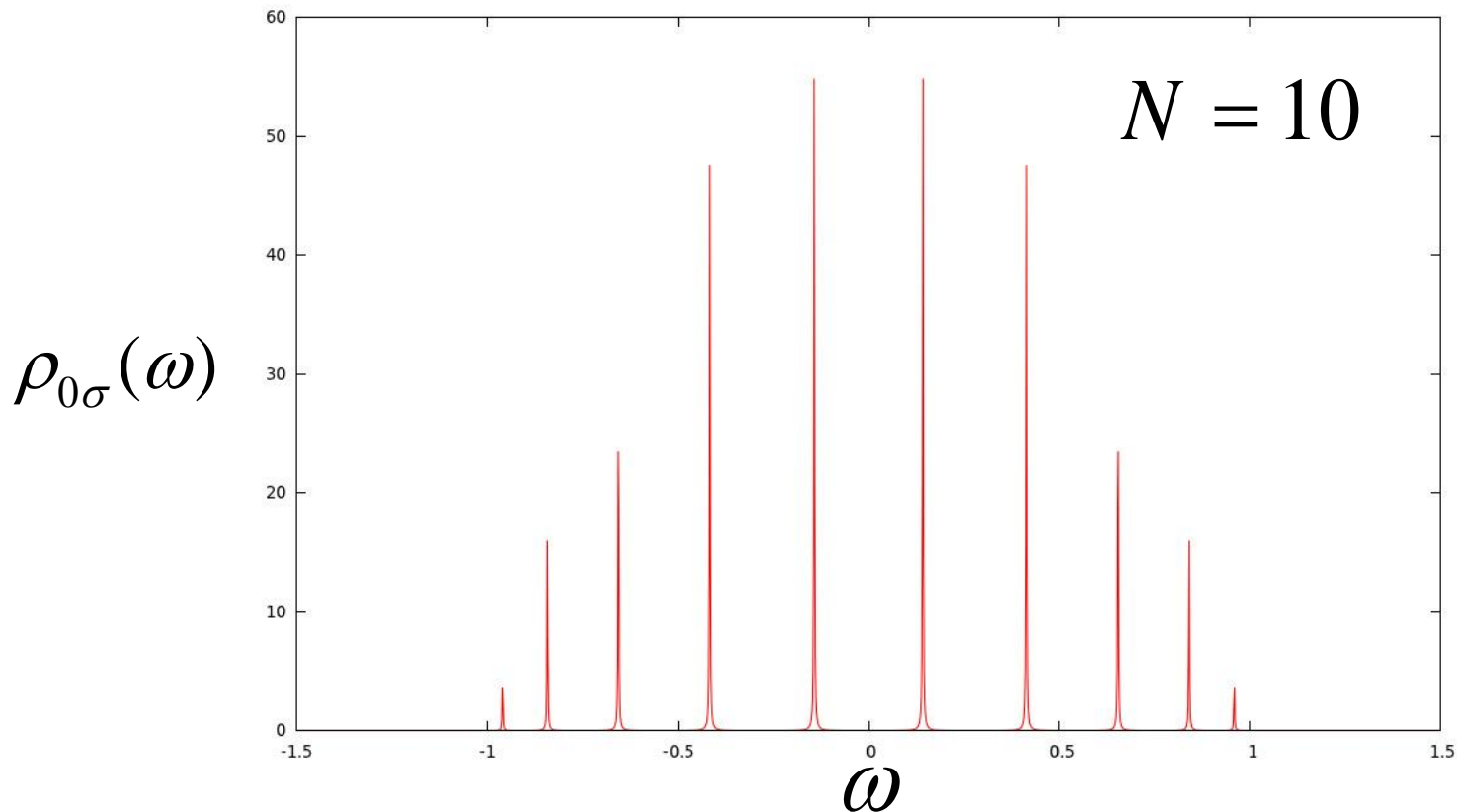
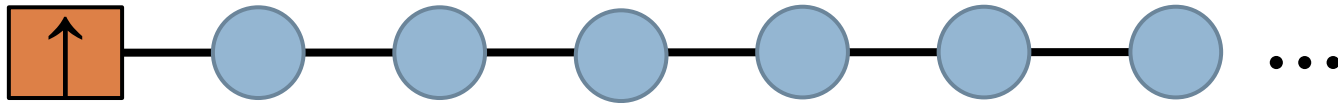
# Impurity problem as a 1d chain



- **Idea: truncate chain?**
  - **Represent bath by the first  $N$  sites of the chain?**
  - **Then exact diagonalization of approximate model?**

# Impurity problem as a 1d chain

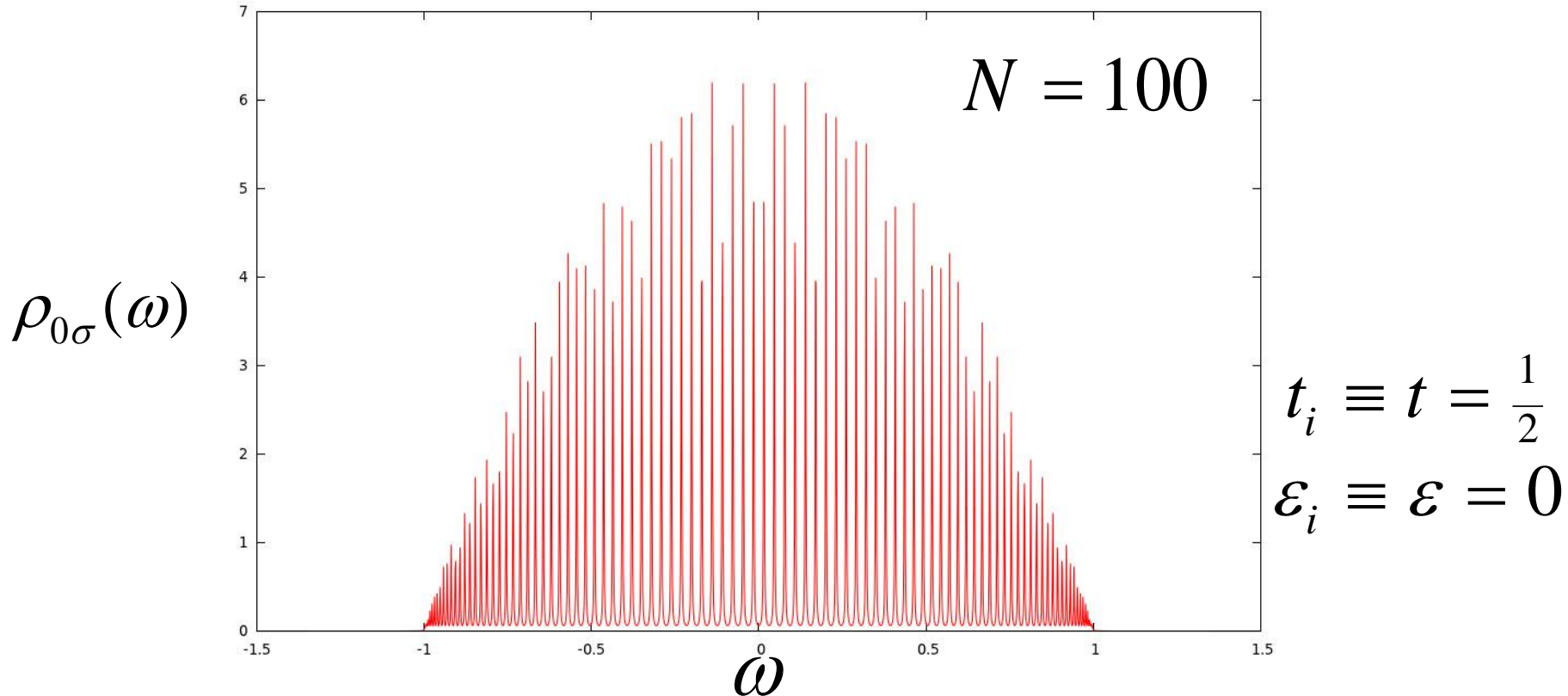
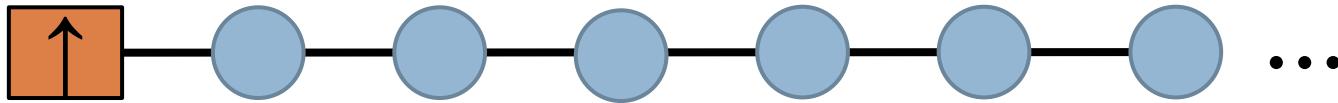
**Bath: 1d chain with constant hopping and onsite energies**



$$t_i \equiv t = \frac{1}{2}$$
$$\varepsilon_i \equiv \varepsilon = 0$$

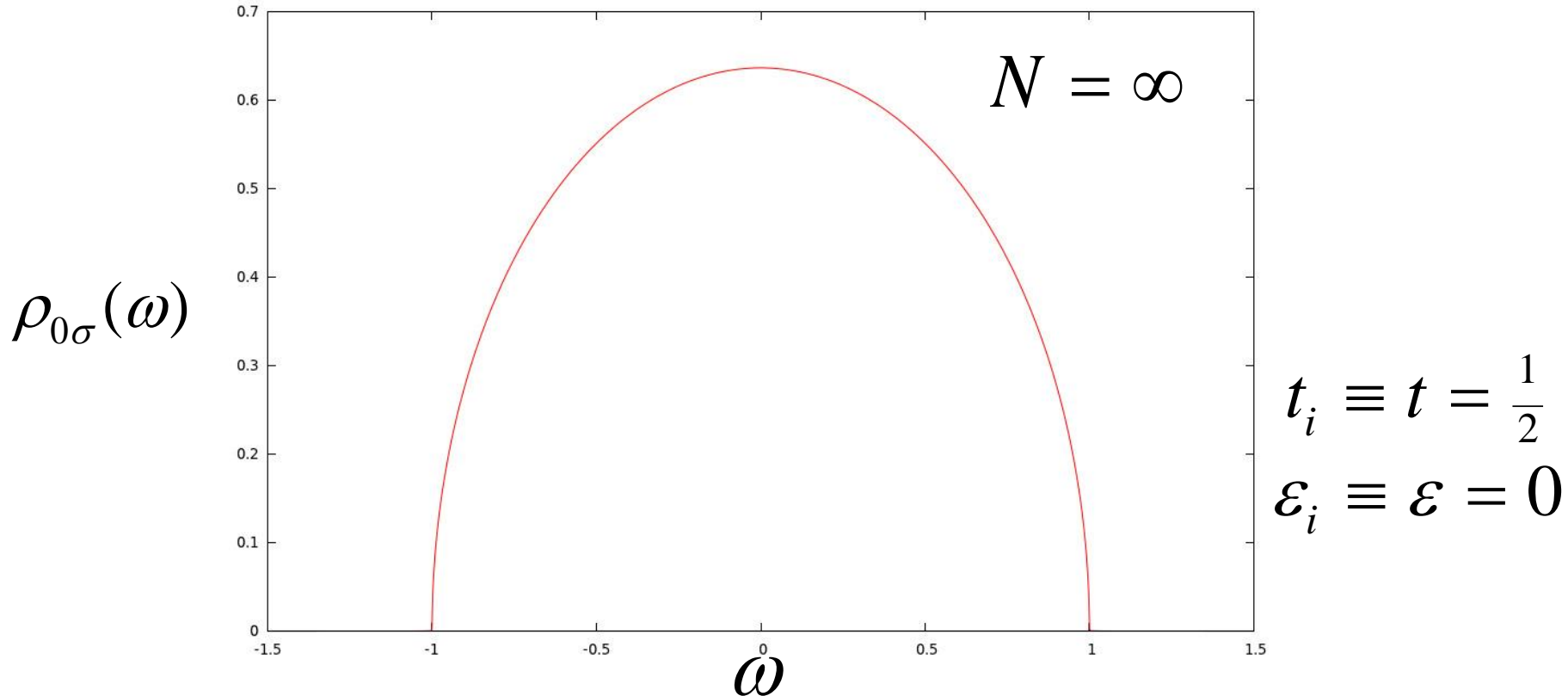
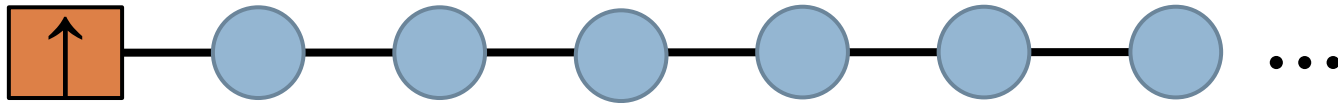
# Impurity problem as a 1d chain

**Bath: 1d chain with constant hopping and onsite energies**

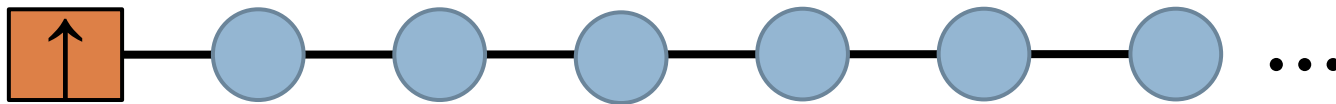


# Impurity problem as a 1d chain

**Bath: 1d chain with constant hopping and onsite energies**



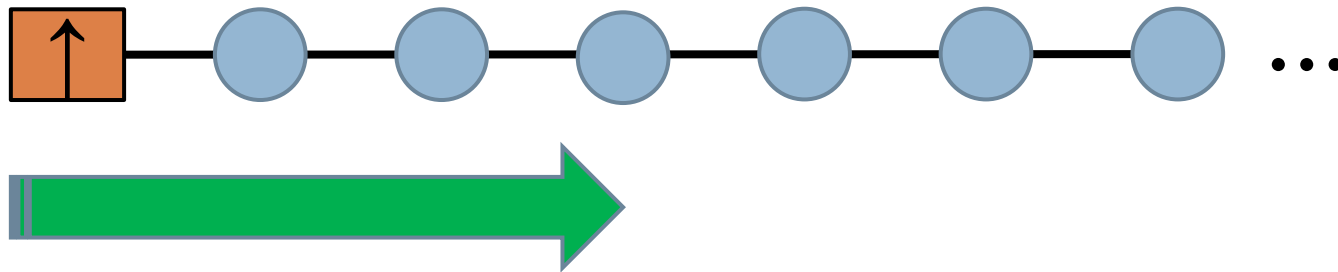
# Impurity problem as a 1d chain



## □ Problem!

- Continuum limit not well-described by finite chain!
- Spectrum has  $N$  poles for an  $N$ -site chain
- Lowest energy scale resolved is of order  $\sim t / N$
- Unable to capture low-energy excitations around Fermi level, which are central to the Kondo effect

# Impurity problem as a 1d chain



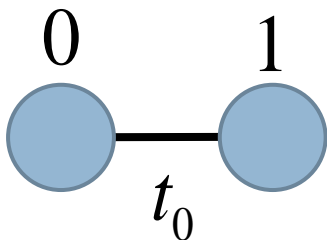
- Idea in NRG: diagonalize the chain **iteratively**
  - Throw away **unimportant** states at each step (successively truncate the Hilbert space)
  - Which states are unimportant?!
    - ... more on that in next lecture!

# Iterative diagonalization

- Take the ‘generic’ tight-binding 1d chain

$$H_N = \sum_{i=0}^{N-1} \sum_{\sigma} t_i c_{i\sigma}^{\dagger} c_{(i+1)\sigma} + \text{H.c.} + \varepsilon_i c_{i\sigma}^{\dagger} c_{i\sigma}$$

- Strategy: build up chain *successively* by adding on extra sites. Start by diagonalizing the dimer:



$$H_1 = \sum_{\sigma} t_0 c_{0\sigma}^{\dagger} c_{1\sigma} + \text{H.c.}$$

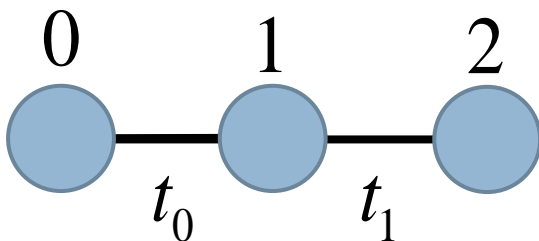


# Iterative diagonalization

- Take the ‘generic’ tight-binding 1d chain

$$H_N = \sum_{i=0}^{N-1} \sum_{\sigma} t_i c_{i\sigma}^{\dagger} c_{(i+1)\sigma} + \text{H.c.}$$

- Strategy: build up chain *successively* by adding on extra sites. Start by diagonalizing the dimer:



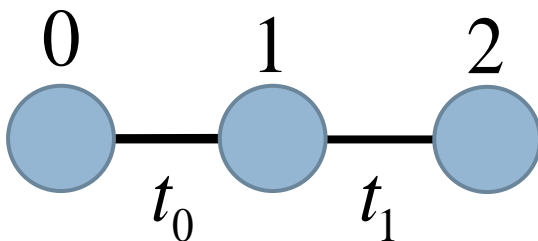
$$H_2 = \sum_{\sigma} t_0 c_{0\sigma}^{\dagger} c_{1\sigma} + t_1 c_{1\sigma}^{\dagger} c_{2\sigma} + \text{H.c.}$$

# Iterative diagonalization

- Take the ‘generic’ tight-binding 1d chain

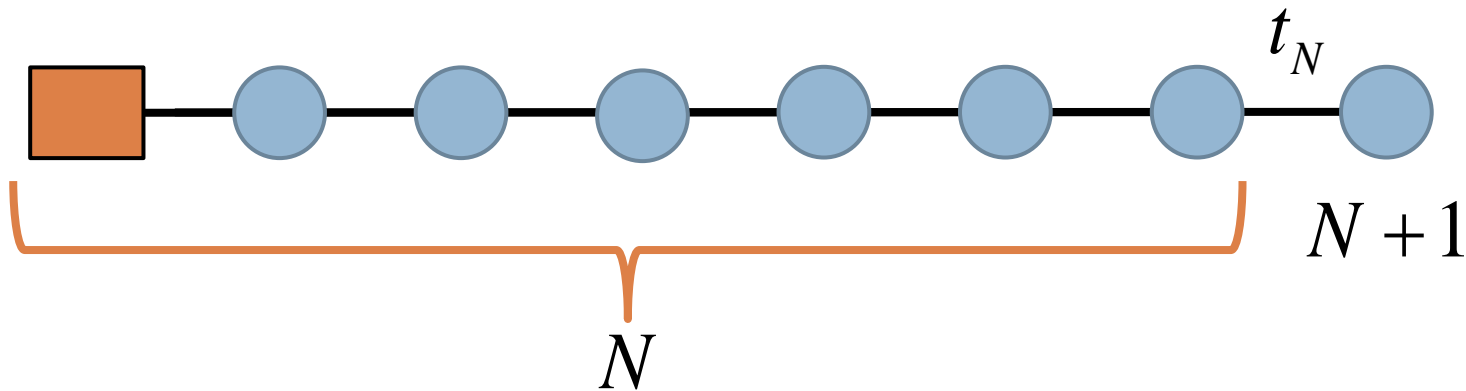
$$H_N = \sum_{i=0}^{N-1} \sum_{\sigma} t_i c_{i\sigma}^{\dagger} c_{(i+1)\sigma} + \text{H.c.}$$

- Strategy: build up chain *successively* by adding on extra sites. Start by diagonalizing the dimer:



$$H_2 = H_1 + \sum_{\sigma} t_1 c_{1\sigma}^{\dagger} c_{2\sigma} + \text{H.c.}$$

# Iterative diagonalization



Recursion:

$$H_{N+1} = H_N + H_{N+1}^{hop}$$

$$H_{N+1}^{hop} = \sum_{\sigma} t_N c_{N\sigma}^{\dagger} c_{(N+1)\sigma} + \text{H.c.}$$

... same transformation for any  $N$  !

# Iterative diagonalization

- Use the **diagonal basis** of iteration  $N$  and couple on an extra site. Then re-diagonalize.
- First, define matrix product states:

$$|N + 1; k, r\rangle_{bs} = |N + 1; k\rangle_a \otimes |N; r\rangle_d$$

**Basis states for iteration N+1**

**Basis states for a single added site**

**Direct product**

**Diagonal states of previous iteration, N**

# Iterative diagonalization

- States of added site defined by:

$$|N + 1; k = 0\rangle_a = |-\rangle = |vac\rangle$$

$$|N + 1; k = +1\rangle_a = |\uparrow\rangle = c_{(N+1)\uparrow}^\dagger |vac\rangle$$

$$|N + 1; k = -1\rangle_a = |\downarrow\rangle = c_{(N+1)\downarrow}^\dagger |vac\rangle$$

$$|N + 1; k = 2\rangle_a = |\uparrow\downarrow\rangle = c_{(N+1)\uparrow}^\dagger c_{(N+1)\downarrow}^\dagger |vac\rangle$$

# Iterative diagonalization

- Diagonalized states expressed as a linear combination of basis states:

$$|N + 1; r\rangle_d = \sum_{k, r'} U_{N+1}^r(k, r') |N + 1; k, r'\rangle_{bs}$$



**Coefficients obtained  
by diagonalizing  
matrix  $H_{N+1}$**

# Iterative diagonalization

□ Construct Hamiltonian:  $H_{N+1} = H_N + H_{N+1}^{hop}$

□ Matrix elements:

$${}_{bs} \langle N+1; k, r | H_{N+1} | N+1; k', r' \rangle_{bs}$$

$$= {}_d \langle N; r | {}_a \langle N+1; k | H_{N+1} | N+1; k' \rangle_a | N; r' \rangle_d$$


$$= {}_d \langle N; r | {}_a \langle N+1; k | H_N | N+1; k' \rangle_a | N; r' \rangle_d$$

$$+ t_N \sum_{\sigma} {}_d \langle N; r | {}_a \langle N+1; k | c_{N\sigma}^{\dagger} c_{(N+1)\sigma} | N+1; k' \rangle_a | N; r' \rangle_d$$

# Iterative diagonalization

$$\begin{aligned} & {}_d \langle N; r | {}_a \langle N + 1; k | H_N | N + 1; k' \rangle_a | N; r' \rangle_d \\ &= {}_a \langle N + 1; k | N + 1; k' \rangle_a \times {}_d \langle N; r | H_N | N; r' \rangle_d \end{aligned}$$

$$\delta_{kk'}$$


$$E_N(r) \delta_{rr'}$$




# Iterative diagonalization

$$t_N \sum_{\sigma} \langle N; r |_a \langle N+1; k | c_{N\sigma}^{\dagger} \hat{\mathbb{1}} c_{(N+1)\sigma} | N+1; k' \rangle_a | N; r' \rangle_d$$



**insert complete set of basis states:**

$$\hat{\mathbb{1}} = \sum_{k'', r''} | N+1; k'' \rangle_a | N; r'' \rangle_d \langle N; r'' |_a \langle N+1; k'' |$$

# Iterative diagonalization

$$(-1)^{k''} \langle N+1; k | N+1; k'' \rangle_a c_{N\sigma}^\dagger$$

$$t_N \sum_{\sigma} \sum_{k'', r''} {}_d \langle N; r | {}_a \langle N+1; k | c_{N\sigma}^\dagger | N+1; k'' \rangle_a | N; r'' \rangle_d$$

$$\times {}_d \langle N; r'' | {}_a \langle N+1; k'' | c_{(N+1)\sigma} | N+1; k' \rangle_a | N; r' \rangle_d$$

# Iterative diagonalization

$$\begin{aligned}
 & \delta_{k,k''} (-1)^k c_{N\sigma}^\dagger \\
 & \underbrace{\hspace{10em}} \\
 t_N \sum_{\sigma} \sum_{k'', r''} & \langle N; r | \langle N+1; k | c_{N\sigma}^\dagger | N+1; k'' \rangle_a | N; r'' \rangle_d \\
 & \times \langle N; r'' | \langle N+1; k'' | c_{(N+1)\sigma} | N+1; k' \rangle_a | N; r' \rangle_d
 \end{aligned}$$

# Iterative diagonalization

$$\delta_{k,k''} (-1)^k \times {}_d \langle N; r | c_{N\sigma}^\dagger | N; r'' \rangle_d$$

$$t_N \sum_{\sigma} \sum_{k'', r''} {}_d \langle N; r | {}_a \langle N+1; k | c_{N\sigma}^\dagger | N+1; k'' \rangle_a | N; r'' \rangle_d$$

$$\times {}_d \langle N; r'' | {}_a \langle N+1; k'' | c_{(N+1)\sigma} | N+1; k' \rangle_a | N; r' \rangle_d$$

# Iterative diagonalization

$$\delta_{k,k''} (-1)^k \times {}_d \langle N; r | c_{N\sigma}^\dagger | N; r'' \rangle_d$$

$$t_N \sum_{\sigma} \sum_{k'', r''} {}_d \langle N; r | {}_a \langle N+1; k | c_{N\sigma}^\dagger | N+1; k'' \rangle_a | N; r'' \rangle_d$$

$$\times {}_d \langle N; r'' | {}_a \langle N+1; k'' | c_{(N+1)\sigma} | N+1; k' \rangle_a | N; r' \rangle_d$$

$$M_{k'', k'}^{\sigma}$$

# Iterative diagonalization

$$\delta_{k,k''} (-1)^k \times {}_d \langle N; r | c_{N\sigma}^\dagger | N; r'' \rangle_d$$

$$\begin{aligned}
 & t_N \sum_{\sigma} \sum_{k'', r''} {}_d \langle N; r | {}_a \langle N+1; k | c_{N\sigma}^\dagger | N+1; k'' \rangle_a | N; r'' \rangle_d \\
 & \times {}_d \langle N; r'' | {}_a \langle N+1; k'' | c_{(N+1)\sigma} | N+1; k' \rangle_a | N; r' \rangle_d \\
 & \qquad \qquad \qquad {}_d \langle N; r'' | N; r' \rangle_d M_{k'', k'}^{\sigma}
 \end{aligned}$$

# Iterative diagonalization

$$\delta_{k,k''} (-1)^k \times \underbrace{{}_d \langle N; r | c_{N\sigma}^\dagger | N; r'' \rangle_d}$$

$$t_N \sum_{\sigma} \sum_{k'', r''} \underbrace{{}_d \langle N; r | {}_a \langle N+1; k | c_{N\sigma}^\dagger | N+1; k'' \rangle_a | N; r'' \rangle_d}_{\delta_{r'', r'} M_{k'', k'}^\sigma} \\ \times \underbrace{{}_d \langle N; r'' | {}_a \langle N+1; k'' | c_{(N+1)\sigma} | N+1; k' \rangle_a | N; r' \rangle_d}$$

# Iterative diagonalization

- Putting it all together:

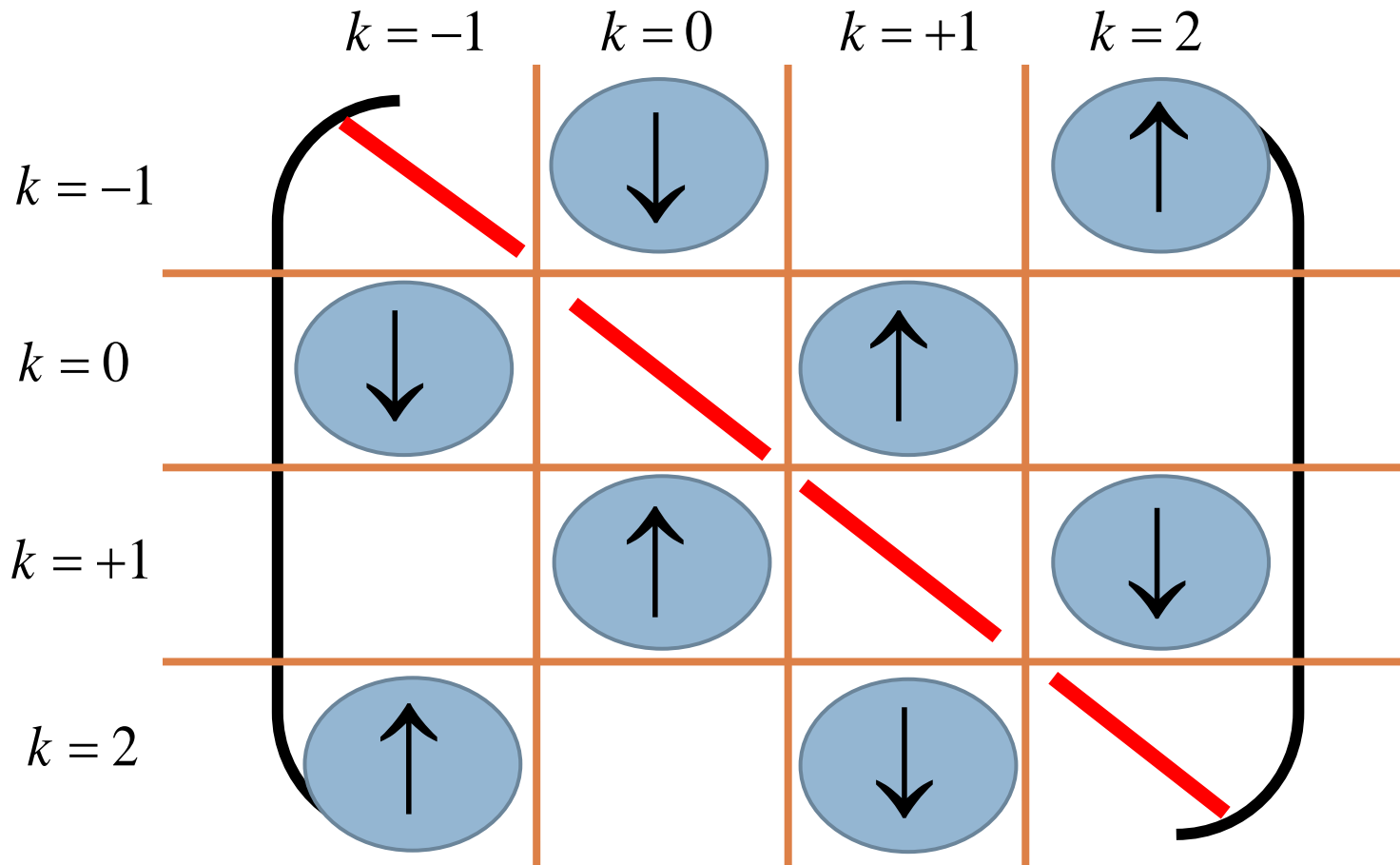
$$\begin{aligned} & {}_{bs} \langle N+1; k, r | H_{N+1} | N+1; k', r' \rangle_{bs} \\ = & E_N(r) \delta_{rr'} \delta_{kk'} + t_N \sum_{\sigma} (-1)^k M_{k,k'}^{\sigma} \times {}_d \langle N; r | c_{N\sigma}^{\dagger} | N; r' \rangle_d \end{aligned}$$

- So: iterative diagonalization requires only:
  - Trivial matrix  $M_{k,k'}^{\sigma}$  which is **independent of  $N$**
  - **Eigenenergies** of previous iteration,  $E_N(r)$
  - **Matrix elements**  ${}_d \langle N; r | c_{N\sigma}^{\dagger} | N; r' \rangle_d$   
between diagonal states of previous iteration



# Iterative diagonalization

- Structure of Hamiltonian matrix:



# Iterative diagonalization

□ How to calculate non-trivial matrix elements?

□ First, **diagonalize** Hamiltonian numerically:

$$|N + 1; r\rangle_d = \sum_{\tilde{k}, \tilde{r}} U_{N+1}^r(\tilde{k}, \tilde{r}) |N + 1; \tilde{k}, \tilde{r}\rangle_{bs}$$

□ For the next iteration, we'll need:

$${}_d \langle N + 1; r | c_{(N+1)\sigma}^\dagger | N + 1; r' \rangle_d = \sum_{\substack{\tilde{r}, \tilde{k} \\ \tilde{r}', \tilde{k}'}} [U_{N+1}^r(\tilde{k}, \tilde{r})]^\dagger U_{N+1}^{r'}(\tilde{k}', \tilde{r}') \times {}_{bs} \langle N + 1; \tilde{k}, \tilde{r} | c_{(N+1)\sigma}^\dagger | N + 1; \tilde{k}', \tilde{r}' \rangle_{bs}$$

# Iterative diagonalization

$$\begin{aligned}
 & {}_{bs} \left\langle N + 1; \tilde{k}, \tilde{r} \left| c_{(N+1)\sigma}^\dagger \right| N + 1; \tilde{k}', \tilde{r}' \right\rangle_{bs} \\
 = & {}_d \left\langle N; \tilde{r} \left| {}_a \left\langle N + 1; \tilde{k} \left| c_{(N+1)\sigma}^\dagger \right| N + 1; \tilde{k}' \right\rangle_a \right| N; \tilde{r}' \right\rangle_d \\
 & \underbrace{\hspace{15em}} \\
 & \left[ M_{\tilde{k}, \tilde{k}'}^\sigma \right]^\dagger \\
 & \delta_{\tilde{r}, \tilde{r}'}
 \end{aligned}$$

# Iterative diagonalization

$${}_d \langle N + 1; r | c_{(N+1)\sigma}^\dagger | N + 1; r' \rangle_d = \sum_{\tilde{r}, \tilde{k}, \tilde{k}'} M_{\tilde{k}', \tilde{k}}^\sigma \left[ U_{N+1}^r(\tilde{k}, \tilde{r}) \right]^\dagger U_{N+1}^{r'}(\tilde{k}', \tilde{r})$$

- **Starting at one end of the chain, we can couple on extra sites recursively, and iteratively diagonalize.**

# Truncation

- **Obvious problem:** Hilbert space grows by a factor of 4 at each iteration (fermionic sites)
  - ▣ After only 10 sites have been added, Hamiltonian matrix in the many-particle basis is of dimension  $10^6 \times 10^6$
  - ▣ Diagonalization time scales as cube of matrix dimension
  - ▣ **Disaster!**
- **Must stop after only a few steps:  
cannot access low-energy physics this way!**

# Truncation

- **Exploit RG concept**  
(which we know is at the heart of quantum impurity problems, from perturbative scaling)
- **Idea in NRG:** throw away **high-energy states** at each iteration, focusing on the low-energy physics at each step. Eventually determine ground state.
- **How do to this?**  
... solution next lecture!